

Chair for Continuum Mechanics Institute of Engineering Mechanics (Prof. Böhlke) Department of Mechanical Engineering

S. Wulfinghoff, T. Böhlke, E. Bayerschen Single Crystal Gradient Plasticity – Part II

Chair for Continuum Mechanics – Institute of Engineering Mechanics



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Literature



- Fundamental dislocation theory
 e.g. Taylor (1934); Orowan (1935); Schmid and Boas (1935); Hall (1951); Petch (1953)
- Kinematics and crystallographic aspects of GNDs Nye (1953); Bilby, Bullough and Smith (1955); Kröner (1958); Mura (1963); Arsenlis and Parks (1999)
- Thermodynamic gradient theories

e.g. Fleck et al. (1994); Steinmann (1996); Menzel and Steinmann (2001); Liebe and Steinmann (2001); Reese and Svendsen (2003); Berdichevski (2006); Ekh et al. (2007); Gurtin, Anand and Lele (2007); Fleck and Willis (2009); Bargmann et al. (2010); Miehe (2011)

- Slip resistance dependent on GNDs and SSDs
 e.g. Becker and Miehe (2004); Evers, Brekelmans and Geers (2004); Cheong, Busso and Arsenlis (2005)
- Micromorphic approach

e.g. Forest (2009); Cordero et al. (2010); Aslan et al. (2011)

Continuum Dislocation Dynamics

e.g. Hochrainer (2006); Hochrainer, Zaiser and Gumbsch (2007); Hochrainer, Zaiser and Gumbsch (2010); Sandfeld, Hochrainer and Zaiser (2010); Sandfeld (2010)

Outline



Motivation

Equivalent Plastic Strain Gradient Plasticity

Numerical results

Dislocation density model

Conclusion and Outlook



Motivation



Typical crystal gradient plasticity extension:

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W(\ldots) \to W(\ldots) + W_g(\nabla \gamma_1, \ \nabla \gamma_2, \ \ldots, \ \nabla \gamma_N)
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- Problem: 15 node variables (FCC) →Gradients massively increase DOFs
- Proposition:

 $W(...) \to W(...) + W_g(\nabla \gamma_{eq})$

- Simplification $\alpha \rightarrow \nabla \gamma_{eq}$
- node variables: $15 \rightarrow 4$



Small Strains

$$oldsymbol{H} =
abla oldsymbol{u} = oldsymbol{H}^e + oldsymbol{H}^p; \hspace{1em} oldsymbol{H}^p = \sum\limits_lpha \lambda_lpha oldsymbol{d}_lpha \otimes oldsymbol{n}_lpha; \hspace{1em} oldsymbol{arepsilon}^p = \sum\limits_lpha \lambda_lpha oldsymbol{M}^S_lpha$$

Two slip parameters λ_{α} per slip system, $\dot{\lambda}_{\alpha} \geq 0$.

- Equivalent plastic strain $\gamma_{eq} \rightarrow \sum_{\alpha} \lambda_{\alpha} = \Sigma_{\lambda}$
- Equality of γ_{eq} and Σ_{λ} is weakly enforced
- Stored Energy $W = W_e(\varepsilon, \varepsilon^p) + W_h(\Sigma_\lambda) + W_g(\nabla \gamma_{eq}) + \check{p}(\Sigma_\lambda \gamma_{eq})$ \check{p} : Lagrange multiplier
- Independent variables: $\boldsymbol{u}, \lambda_{\alpha}, \gamma_{eq}$



Variation of displacements ($\delta \lambda_{\alpha} = \delta \gamma_{eq} = 0$)

$$\delta \int_{\Delta \mathcal{B}} W \, \mathrm{d}v = \int_{\Delta \mathcal{B}} \partial_{\boldsymbol{\varepsilon}} W \cdot \delta \boldsymbol{\varepsilon} \, \mathrm{d}v = \int_{\partial \Delta \mathcal{B}} \boldsymbol{t} \cdot \delta \boldsymbol{u} \, \mathrm{d}a \tag{1}$$

$$\Leftrightarrow \int_{\partial \Delta \mathcal{B}} (\partial_{\varepsilon} W \boldsymbol{n} - \boldsymbol{t}) \cdot \delta \boldsymbol{u} \, \mathrm{d}\boldsymbol{a} - \int_{\Delta \mathcal{B}} \operatorname{div} (\partial_{\varepsilon} W) \cdot \delta \boldsymbol{u} \, \mathrm{d}\boldsymbol{v} = 0$$
(2)

Result:

- Cauchy stress $\sigma = \partial_{\boldsymbol{\varepsilon}} W$
- Linear momentum balance $\operatorname{div}(\boldsymbol{\sigma}) = 0$





Variation of γ_{eq} ($\delta \lambda_{\alpha} = 0, \ \delta u = 0$)

$$\begin{split} \delta \int_{\Delta \mathcal{B}} W \, \mathrm{d}v &= \int_{\Delta \mathcal{B}} (\underbrace{\partial_{\gamma_{eq}} W}_{-\check{p}} \, \delta\gamma_{eq} + \underbrace{\partial_{\nabla\gamma_{eq}} W}_{\boldsymbol{\xi}} \cdot \nabla(\delta\gamma_{eq})) \, \mathrm{d}v = \int_{\partial \Delta \mathcal{B}} \Xi \, \delta\gamma_{eq} \, \mathrm{d}a \\ \Leftrightarrow \int_{\partial \Delta \mathcal{B}} (\boldsymbol{\xi} \cdot \boldsymbol{n} - \Xi) \cdot \delta \boldsymbol{u} \, \mathrm{d}a - \int_{\Delta \mathcal{B}} (\operatorname{div}(\boldsymbol{\xi}) + \check{p}) \delta\gamma_{eq} \, \mathrm{d}v = 0 \end{split}$$

Result:

- Micro traction $\Xi = \boldsymbol{\xi} \cdot \boldsymbol{n}$
- Backstress $\check{p} = -\operatorname{div}(\boldsymbol{\xi})$





Dissipation and flow rule

$$D_{tot}(\Delta \mathcal{B}) = \int_{\partial \Delta \mathcal{B}} (\boldsymbol{t} \cdot \dot{\boldsymbol{u}} + \Xi \, \dot{\gamma}_{eq}) \, \mathrm{d}\boldsymbol{a} - \int_{\Delta \mathcal{B}} \dot{W} \, \mathrm{d}\boldsymbol{v} \stackrel{!}{\geq} 0 \tag{3}$$

with $\operatorname{div}(\sigma) = 0$, $\sigma n = t$, $\Xi = \boldsymbol{\xi} \cdot \boldsymbol{n}$ and $\check{p} = -\operatorname{div}(\boldsymbol{\xi})$:

$$D_{tot}(\Delta \mathcal{B}) = \int_{\Delta \mathcal{B}} \underbrace{\sum_{\alpha} (\boldsymbol{\sigma} \cdot \boldsymbol{M}_{\alpha}^{S} - \partial_{\Sigma_{\lambda}} W_{h} - \check{p}) \dot{\lambda}_{\alpha}}_{\mathcal{D}} dv \stackrel{!}{\geq} 0$$
(4)

Local dissipation $\mathcal{D} \stackrel{!}{\geq} 0 \Rightarrow$ Flow rule, e.g. :

$$\dot{\lambda}_{\alpha} = \dot{\gamma}_0 \left\langle \frac{\tau_{\alpha} - (\tau_0^C + \partial_{\Sigma_{\lambda}} W_h + \check{p})}{\tau^D} \right\rangle^p$$

(5)



Global residuals for FE-Implementation

$$G^{u} = \int_{\mathcal{B}} \boldsymbol{\sigma} \cdot \delta \boldsymbol{\varepsilon} \, \mathrm{d}\boldsymbol{v} - \int_{\partial \mathcal{B}_{t}} \bar{\boldsymbol{t}} \cdot \delta \boldsymbol{u} \, \mathrm{d}\boldsymbol{a} \stackrel{!}{=} 0$$
$$G^{\gamma_{eq}} = \int_{\mathcal{B}} (\boldsymbol{\xi} \cdot \nabla(\delta \gamma_{eq}) - \check{p} \, \delta \gamma_{eq}) \, \mathrm{d}\boldsymbol{v} - \int_{\partial \mathcal{B}_{\Xi}} \bar{\Xi} \, \delta \gamma_{eq} \, \mathrm{d}\boldsymbol{a} \stackrel{!}{=} 0$$

Local (integration point) residuals based on implicit Euler:

$$\boldsymbol{r}^{\sigma} = \underbrace{\mathbb{C}^{-1}[\boldsymbol{\sigma}]}_{\boldsymbol{\varepsilon}^{e}} - \left(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_{n}^{p} - \sum_{\alpha} \Delta \lambda_{\alpha}(\boldsymbol{\sigma}, \boldsymbol{p}) \boldsymbol{M}_{\alpha}^{S}\right) \stackrel{!}{=} \boldsymbol{0}$$
$$r^{p} = \gamma_{eq} - \left(\gamma_{eq,n} + \sum_{\alpha} \Delta \lambda_{\alpha}(\boldsymbol{\sigma}, \boldsymbol{p})\right) \stackrel{!}{=} \boldsymbol{0}$$

Possible Enhancements/Modifications



Penalty approximation

$$W = W_e(\varepsilon, \varepsilon^p) + W_h(\Sigma_\lambda) + W_g(\nabla \gamma_{eq}) + \check{p}(\Sigma_\lambda - \gamma_{eq})$$

$$\rightarrow \qquad W = W_e(\varepsilon, \varepsilon^p) + W_h(\Sigma_\lambda) + W_g(\nabla \gamma_{eq}) + \frac{1}{2} H_{\chi}(\Sigma_\lambda - \gamma_{eq})^2$$

Penalty parameter H_χ: large number
 Energy 1/2H_χ(Σ_λ − γ_{eq})² penalizes deviations of Σ_λ from γ_{eq}

Grain Boundary Yield Criterion

$$f_{\Gamma} = \llbracket \Xi \rrbracket - \Xi^C$$

- Ξ^C : Grain boundary yield stress
- Grain boundary yield resistance
- Delays initiation of plastic flow on the grain boundaries

Simulation Results





Figure: Periodic Simulation, 2 Mio. Dof

Hardening model $W_g = \frac{1}{2} K_G \nabla \gamma_{eq} \cdot \nabla \gamma_{eq}; \quad W_h(\Sigma_\lambda) = \tau_\infty^C \Sigma_\lambda + \frac{1}{\theta_0} (\tau_\infty^C - \tau_0^C)^2 \exp\left(-\frac{\theta_0 \Sigma_\lambda}{\tau_\infty^C - \tau_0^C}\right)$ Material Parameters

E	ν	$\dot{\gamma}_0$	p	τ^D	τ_0^C	τ^C_{∞}	θ_0	K_G
$70 { m GPa}$	0.34	$10^{-3} \mathrm{s}^{-1}$	50	$1 \mathrm{MPa}$	$50 \mathrm{MPa}$	$100 \mathrm{MPa}$	$500 \mathrm{MPa}$	0.01 N

Macro Response





Comparison With Hall-Petch Relation





Figure: Tensile test with periodic 2-grain microstructure

Convergence



$\Delta t : 0.012 \text{ s}$ 0.192 s 0.295 s $0.024 \ s$ 0.048 s 0.096 s 0.333 s 1.16e+0.32 33e+03 4 64e+03 9 19e+03 1 83e+04 3 12e+04 2 66e+04 5.20e-01 5.79e+00 1.15e+04 2.55e+04 4.03e+043.85e+02 3.45e+02911e-04 2 28e-01 8 26e+01 1 73e+02 2 66e+02 1 64e+02 1 87e+02 1 35e-05 2 20e-04 2.35e+01 3 15e+01 4 13e+01 3.84e+01 9.05e+01 Macro response 2.01e-07 3.31e-06 3.71e+00 4.30e+00 6.23e+00 9.57e+00 2.09e+01 750 2 976-09 5 00e-08 2 10e-01 4 24e-01 5 65e-01 4 12e+00 4 87e+00 1.63e-02 2.91e-02 1.88e-02 5.67e-01 1.18e+00 2.48e-04 3.64e-04 5.44e-05 7.97e-03 1.53e-01 4 37e+02 500 5 35e-07 5 21e-07 3 94e-05 2 35e-02 MPa 3.16e-09 6.48e+01 9.35e+01 1.24e-06 1.26e-04 1.22e+00 3.19e+011.03e-06 ıb 5 01e-02 3 00e+00 250 Active Set Search 3.10e-04 2.80e-01 Grain Boundary Nodes 8 44e-07 9 25e-03 5 47e-05 2 13e+01 0 2.46e-01 4.67e+01 2 56e-03 6 73e-01 3.49e-06 8.37e-03 4.66e+001.71e-05 8 48e-03 1 47e-07 1.41e-05 2 20e-01 1 28e-04 1.79e-07 3 56e-09

Euclidean Norm of Force Residual

Marked in red: Active set search of grain boundary nodes

0.025

 $\bar{\varepsilon}$ [-]

0.05

Dependence on Time Step Size





Improvement of Cross Hardening Model



- Cross hardening is modelled by $W_h(\Sigma_{\lambda})$ \rightarrow phenomenological approach
- Replacement by dislocation based approach $\tau^{C} = \tau^{C}(\rho) = \tau^{C}_{0} + aGb\sqrt{\rho}$
- Advanced dislocation density model (until now only single slip)

$$\partial_t \rho = -\operatorname{div} \left(\boldsymbol{\kappa}^{\perp} \boldsymbol{\nu} \right) + \overline{\rho k} \boldsymbol{\nu}$$

$$\partial_t \overline{\rho k} = -\operatorname{div} \left(\overline{\rho k} / \rho \, \boldsymbol{\nu} \boldsymbol{\kappa}^{\perp} + \rho / 2 \, \nabla_{\!\! p} \boldsymbol{\nu} \right)$$

$$\dot{\lambda} = \rho b \boldsymbol{\nu}$$

$$\boldsymbol{\nu} = \nu_0 \operatorname{sgn} \left(\tau \right) \left\langle \frac{|\tau| - \tau^C - \check{p}}{\tau^D} \right\rangle^p$$

with $\kappa^{\perp} = 1/b(I - n \otimes n)\nabla\lambda$ and $\nabla_{p}\nu = (I - n \otimes n)\nabla\nu$. Accounts for dislocation transport and curvature-induced line length production.

Summary



- Crystal plasticity model enhanced by gradient of equivalent plastic strain
- DOFs per node massively reduced
- Leads to backstress formally similar to other theories
- Classical FE-implementation + Grain boundary yield condition
- Converging results
- Results close to Hall Petch relation





End of Part II

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