

S. Wulfinghoff, T. Böhlke, E. Bayerschen

# Single Crystal Gradient Plasticity – Part I

Chair for Continuum Mechanics – Institute of Engineering Mechanics

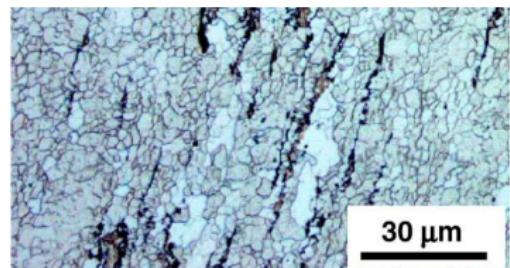
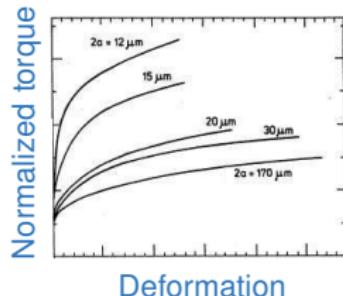
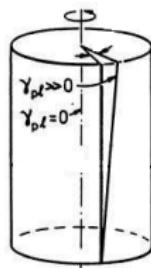
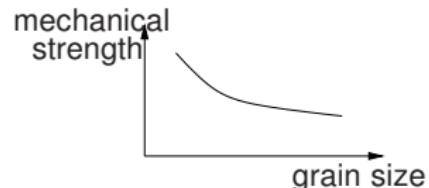


# Motivation

Materials exhibit strong size effects, when the length scale associated with non-uniform plastic deformation is in the order of microns.

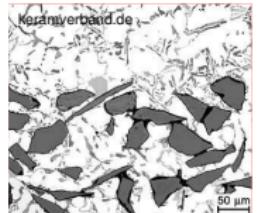
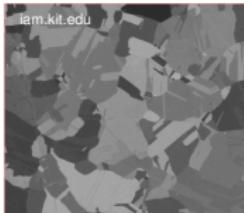
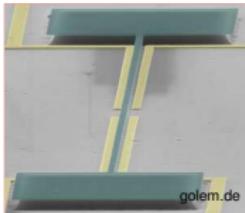
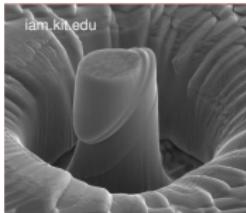
Examples:

- Fine-grained steels
- Precipitation hardening
- Torsion of thin wires [Fleck et al. \(1994\)](#)



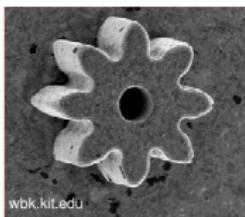
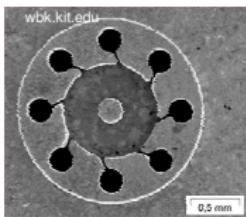
B. Eghbali (2007)

# Motivation



Understanding micro-plasticity

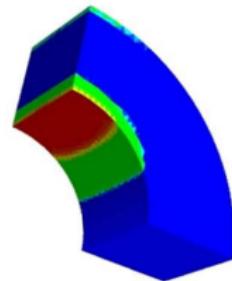
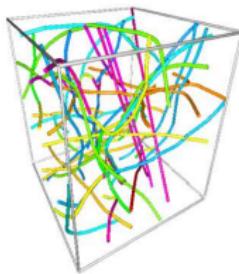
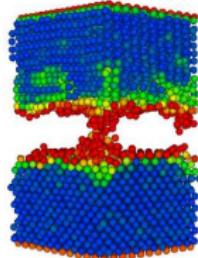
Design of materials



Dimensioning of micro-components and micro-systems

How to identify physically based continuum mechanical micro-plasticity models?

# Motivation

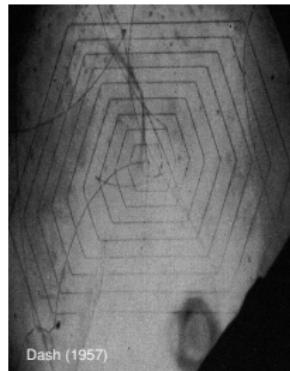


- Atomistic and DDD-simulations are limited to very small systems
- Classical continuum mechanical plasticity fails at the micro-scale

Need for a **Continuum Dislocation Theory** to bridge the scales

DFG Research Group 1650 "Dislocation based Plasticity"

# Dislocation Microstructure



Dash (1957)

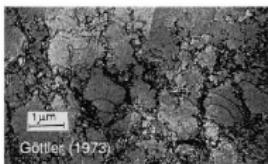


Mughrabi (1971)

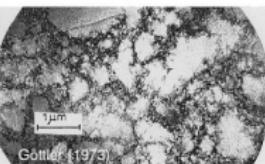


(a)

(b)



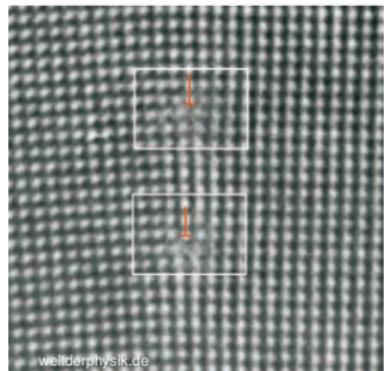
Göttsche (1973)



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Spiral source

Cell structures at different deformation stages



weldphysik.de

Single dislocations

Important features of the microstructure

- Total line length/density
- Dislocation sources
- Dislocation motion/transport
- Lattice distortion

- Fundamental dislocation theory
  - e.g. Taylor (1934); Orowan (1935); Schmid and Boas (1935); Hall (1951); Petch (1953)
- Kinematics and crystallographic aspects of GNDs
  - Nye (1953); Bilby, Bullough and Smith (1955); Kröner (1958); Mura (1963); Arsenlis and Parks (1999)
- Thermodynamic gradient theories
  - e.g. Fleck et al. (1994); Steinmann (1996); Menzel and Steinmann (2001); Liebe and Steinmann (2001); Reese and Svendsen (2003); Berdichevski (2006); Ekh et al. (2007); Gurtin, Anand and Lele (2007); Fleck and Willis (2009); Bargmann et al. (2010); Miehe (2011)
- Slip resistance dependent on GNDs and SSDs
  - e.g. Becker and Miehe (2004); Evers, Brekelmans and Geers (2004); Cheong, Busso and Arsenlis (2005)
- Micromorphic approach
  - e.g. Forest (2009); Cordero et al. (2010); Aslan et al. (2011)
- Continuum Dislocation Dynamics
  - e.g. Hochrainer (2006); Hochrainer, Zaiser and Gumbsch (2007); Hochrainer, Zaiser and Gumbsch (2010); Sandfeld, Hochrainer and Zaiser (2010); Sandfeld (2010)

# Outline

- Single crystal kinematics at small strains
- Free-energy ansatz
- Field and boundary equations
- Constitutive modeling
- Numerical examples

# Kinematics of Single Crystals at Small Strains

- Displacement gradient

$$\mathbf{H} = \mathbf{H}_e + \mathbf{H}_p$$

- Plastic part of the displacement gradient

$$\mathbf{H}_p = \sum_{\alpha} \gamma_{\alpha} \mathbf{d}_{\alpha} \otimes \mathbf{n}^{\alpha}$$

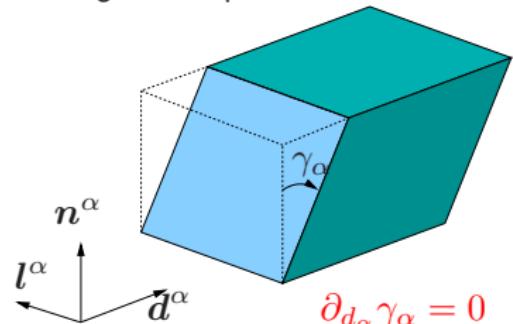
- Additive decomposition of the strain tensor

$$\boldsymbol{\varepsilon} = \text{sym}(\mathbf{H}) = \boldsymbol{\varepsilon}_e + \boldsymbol{\varepsilon}_p$$

- Flow rule

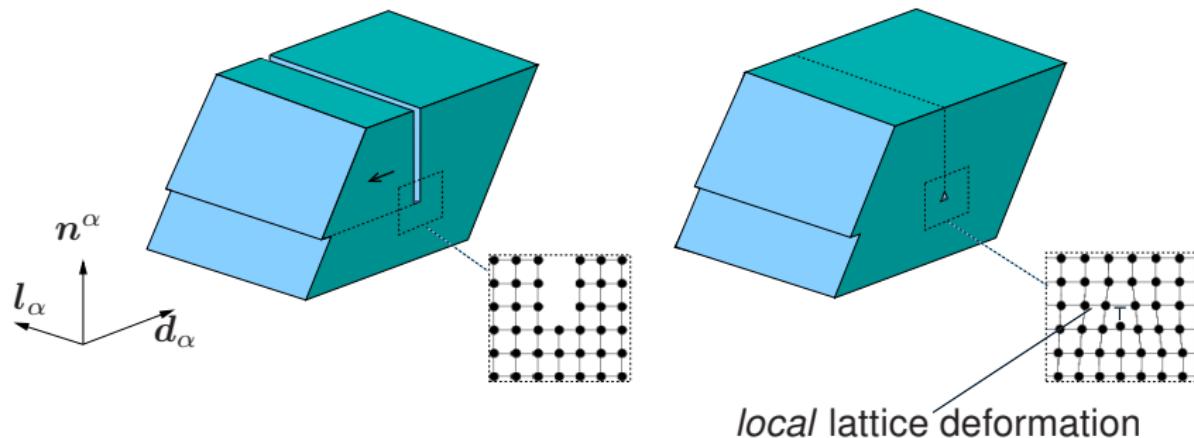
$$\dot{\boldsymbol{\varepsilon}}_p = \sum_{\alpha} \dot{\gamma}_{\alpha} \mathbf{M}_{\alpha} \quad \mathbf{M}_{\alpha} = \frac{1}{2} (\mathbf{d}_{\alpha} \otimes \mathbf{n}^{\alpha} + \mathbf{n}^{\alpha} \otimes \mathbf{d}_{\alpha})$$

Example:  
Homogeneous plastic shear



# Kinematics of Single Crystals at Small Strains

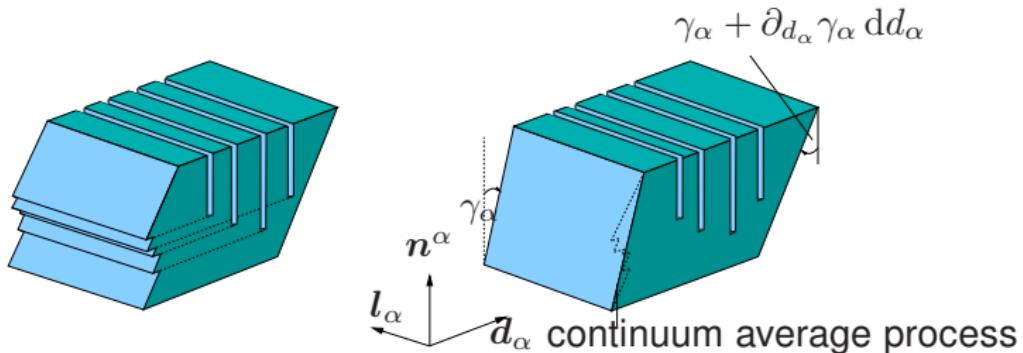
Introduction of a pure edge dislocation (with negative sign)



- Real crystal  $\Rightarrow$  *discontinuous* slip
- *Local* lattice deformation **not** captured precisely by classical theory

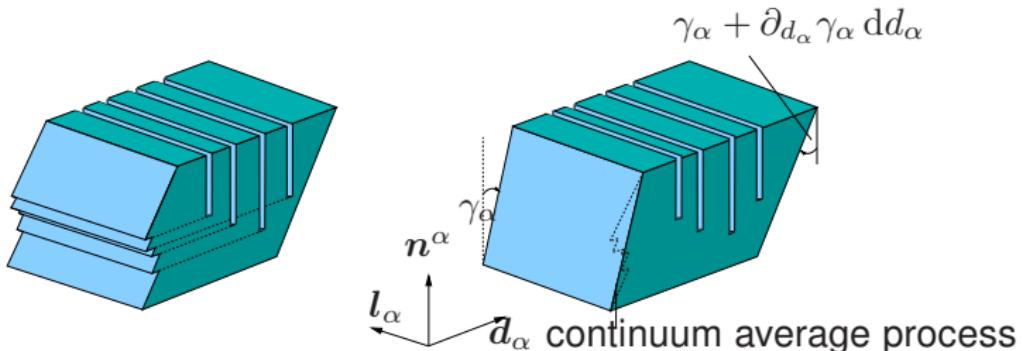
# Kinematics of Single Crystals at Small Strains

Graphical representation of the averaging process of  $\gamma_\alpha$  in continuum mechanics



# Kinematics of Single Crystals at Small Strains

Graphical representation of the averaging process of  $\gamma_\alpha$  in continuum mechanics



⇒ Dislocations of *equal sign* lead to **non-homogeneous** plastic shear

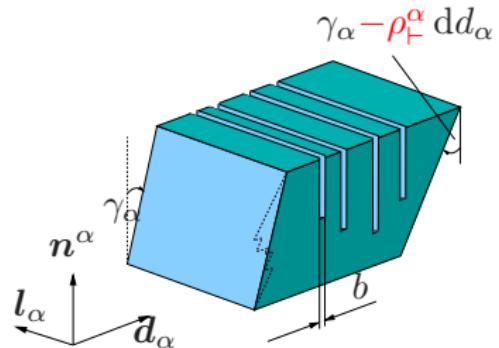
$$\partial_{d_\alpha} \gamma_\alpha \neq 0$$

⇒ *Geometrically necessary* dislocations

# Kinematics of Single Crystals at Small Strains

## Edge and screw dislocation densities

$$\rho_F^\alpha := -\partial_{d_\alpha} \gamma_\alpha = -\nabla \gamma_\alpha \cdot d_\alpha$$



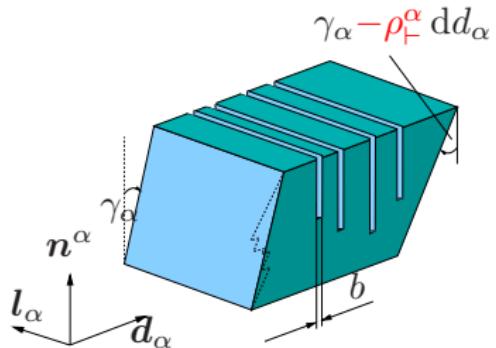
### ■ Dislocation density tensor

$$\alpha = \sum_\alpha [(-\nabla \gamma_\alpha \cdot d_\alpha) l_\alpha \otimes d_\alpha + (\nabla \gamma_\alpha \cdot l_\alpha) d_\alpha \otimes d_\alpha]$$

# Kinematics of Single Crystals at Small Strains

## Edge and screw dislocation densities

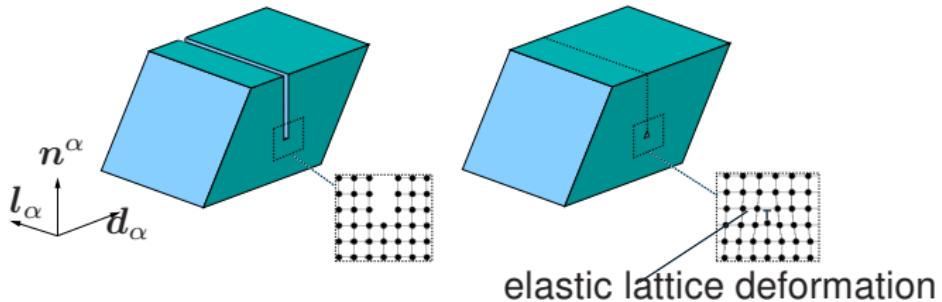
$$\begin{aligned}\rho_F^\alpha &:= -\partial_{d_\alpha} \gamma_\alpha = -\nabla \gamma_\alpha \cdot \mathbf{d}_\alpha \\ \rho_C^\alpha &:= \partial_{l_\alpha} \gamma_\alpha = \nabla \gamma_\alpha \cdot \mathbf{l}_\alpha\end{aligned}$$



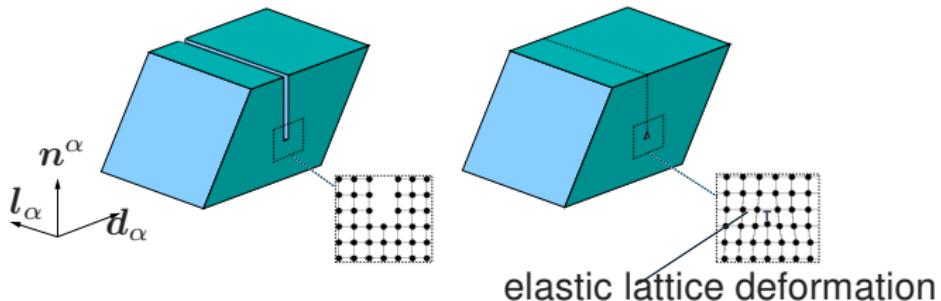
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$$\boldsymbol{\alpha} = \sum_\alpha [(-\nabla \gamma_\alpha \cdot \mathbf{d}_\alpha) \mathbf{l}_\alpha \otimes \mathbf{d}_\alpha + (\nabla \gamma_\alpha \cdot \mathbf{l}_\alpha) \mathbf{d}_\alpha \otimes \mathbf{d}_\alpha]$$

# Free Energy Ansatz

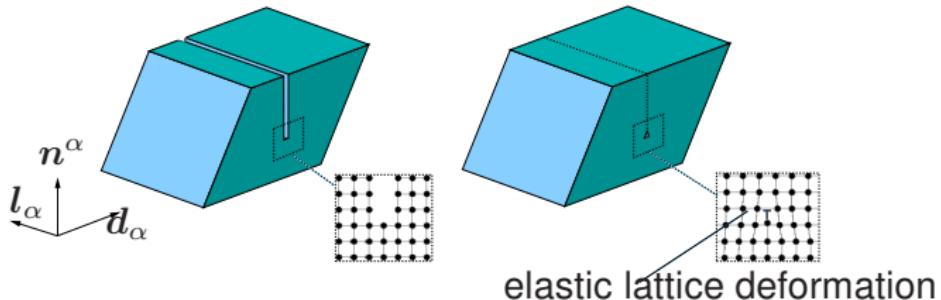


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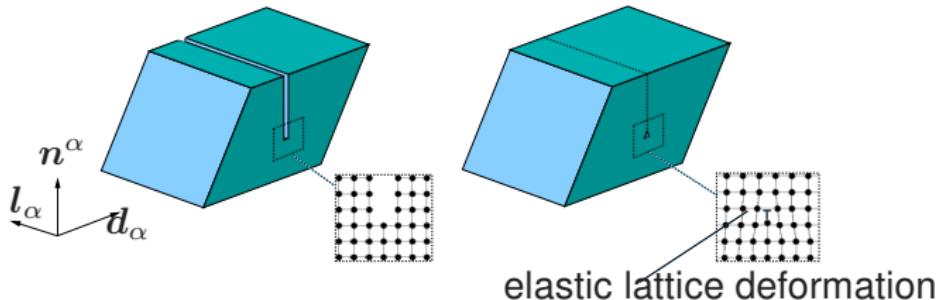


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# Free Energy Ansatz



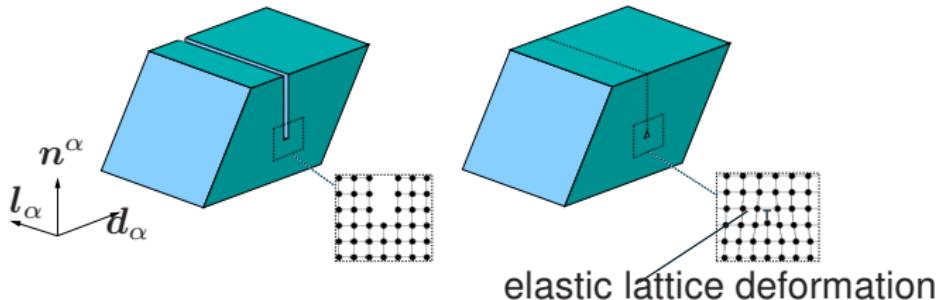
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 $\Rightarrow$  Free energy  $W \geq W_e := \frac{1}{2}\varepsilon_e \cdot \mathbb{C}[\varepsilon_e]$



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## Ansatz

- Free energy  $W = W_e(\varepsilon_e) + W_d(\rho_F^\alpha, \rho_\odot^\alpha)$
- Dissipation  $\mathcal{D}_{tot} = \sum_\alpha \mathcal{D}(\dot{\gamma}_\alpha) = \sum_\alpha \tau_\alpha^D \dot{\gamma}_\alpha$

# Field and Boundary Equations

Virtual work of external forces  $\delta W_{ext}$  and internal forces  $\delta W_{int}$

$$\delta W_{ext} \stackrel{!}{=} \delta W_{int} \quad \forall \delta \mathbf{u}, \delta \gamma_\alpha$$

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$$\delta W_{ext} \stackrel{!}{=} \delta W_{int} \quad \forall \delta \mathbf{u}, \delta \gamma_\alpha$$

$$\begin{aligned}\delta W_{int} &= \delta \int_V W dV + \int_V \sum_{\alpha} \tau_{\alpha}^{\mathcal{D}} \delta \gamma_{\alpha} dV \\ &= \int_V \left[ \partial \boldsymbol{\varepsilon}_e W_e \cdot \delta \boldsymbol{\varepsilon}_e + \sum_{\alpha} \left( \partial \rho_{\vdash}^{\alpha} W_d \delta \rho_{\vdash}^{\alpha} + \partial \rho_{\odot}^{\alpha} W_d \delta \rho_{\odot}^{\alpha} \right) \right] dV \\ &\quad + \int_V \sum_{\alpha} \tau_{\alpha}^{\mathcal{D}} \delta \gamma_{\alpha} dV.\end{aligned}$$

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$$\delta W_e = \boldsymbol{\sigma} \cdot \delta \boldsymbol{\varepsilon}_e$$

$$\delta W_d = \sum_{\alpha} \boldsymbol{\xi}_{\alpha} \cdot \nabla \delta \gamma_{\alpha}$$

$$\boldsymbol{\sigma} := \partial_{\boldsymbol{\varepsilon}_e} W_e$$

$$\boldsymbol{\xi}_{\alpha} := -\partial_{\rho_{\vdash}^{\alpha}} W_d \mathbf{d}_{\alpha} + \partial_{\rho_{\odot}^{\alpha}} W_d \mathbf{l}_{\alpha}$$

stress

defect stress

# Field and Boundary Equations

Principle of virtual displacements yields

## Field equations

$$\operatorname{div}(\boldsymbol{\sigma}) + \mathbf{f} = \mathbf{0}, \quad \tau_{\alpha}^{\mathcal{D}} = \tau_{\alpha} - (-1)\operatorname{div}(\boldsymbol{\xi}_{\alpha}) \quad \forall \mathbf{x} \in V$$

- $\tau_{\alpha} := \boldsymbol{\sigma} \cdot \mathbf{M}_{\alpha}$  resolved shear stress
- $(-1)\operatorname{div}(\boldsymbol{\xi}_{\alpha})$  backstress due to GNDs

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## Boundary equations

$$\boldsymbol{\sigma} \mathbf{n} = \mathbf{t}, \quad \boldsymbol{\xi}_{\alpha} \cdot \mathbf{n} = \Xi_{\alpha} \quad \forall \mathbf{x} \in \partial V_{\sigma}$$

$$\mathbf{u} = \bar{\mathbf{u}}, \quad \gamma_{\alpha} = \bar{\gamma}_{\alpha} \quad \forall \mathbf{x} \in \partial V_u$$

- $\mathbf{t}$  traction
- $\Xi_{\alpha}$  microtraction, results from evaluation of the PovD

## Strain energy (linearized theory)

$$W_e = \frac{1}{2} \boldsymbol{\varepsilon}_e \cdot \mathbb{C}[\boldsymbol{\varepsilon}_e]$$

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## Dissipation

$$\mathcal{D}(\dot{\gamma}_\alpha) = \dot{\gamma}_0 \tau_0 \left| \frac{\dot{\gamma}_\alpha}{\dot{\gamma}_0} \right|^{m+1} \Rightarrow \tau_\alpha^{\mathcal{D}} = \text{sgn}(\dot{\gamma}_\alpha) \tau_0 \left| \frac{\dot{\gamma}_\alpha}{\dot{\gamma}_0} \right|^m$$

## Defect energy (quadratic)

$$W_d = \frac{1}{2} \vec{\rho} \cdot \underline{\underline{E}} \cdot \vec{\rho}, \quad \vec{\rho} = (\rho_{\leftarrow}^1, \rho_{\leftarrow}^2, \dots, \rho_{\leftarrow}^N, \rho_{\odot}^1, \rho_{\odot}^2, \dots, \rho_{\odot}^N)$$

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- Leads to linear dependence of the defect stress on the dislocation densities
- Example – uncoupled quadratic ansatz:

$$W_d = \frac{1}{2} S_0 L^2 \sum_{\alpha} [|\rho_{\leftarrow}^{\alpha}|^2 + |\rho_{\odot}^{\alpha}|^2]; \quad S_0, L = \text{const.}$$

Gurtin et al. (2007)

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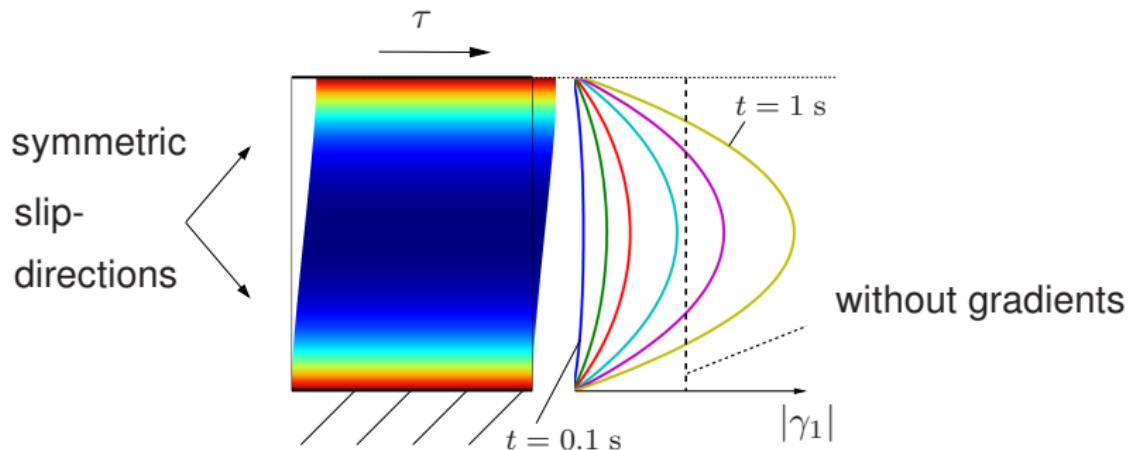
Gurtin et al. (2007)

## Generalized nonlinear ansatz

$$W_d = c_n \sum_{\alpha} [|\rho_{\leftarrow}^{\alpha} L|^n + |\rho_{\odot}^{\alpha} L|^n]; \quad c_n, n, L = \text{const.}$$

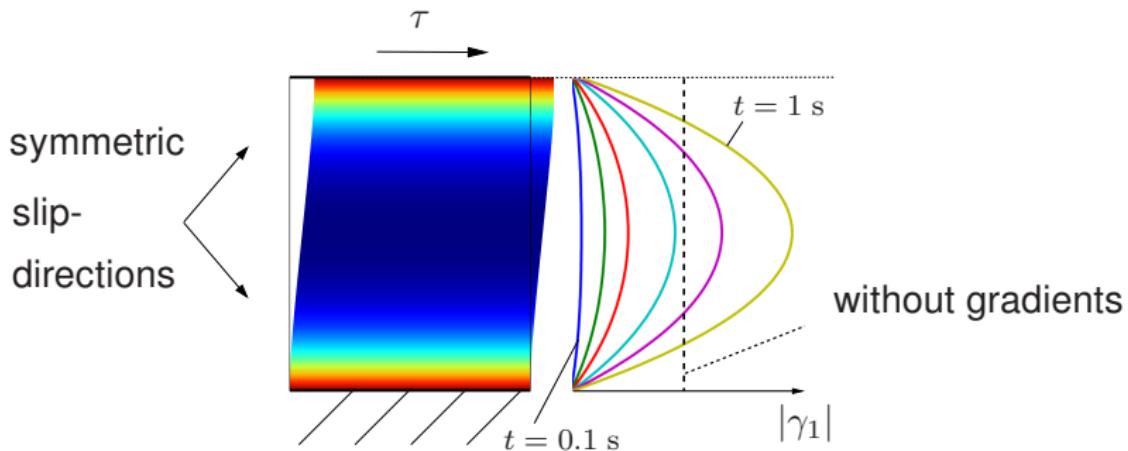
# Numerical Results

Infinite 2d-layer, two slip systems:



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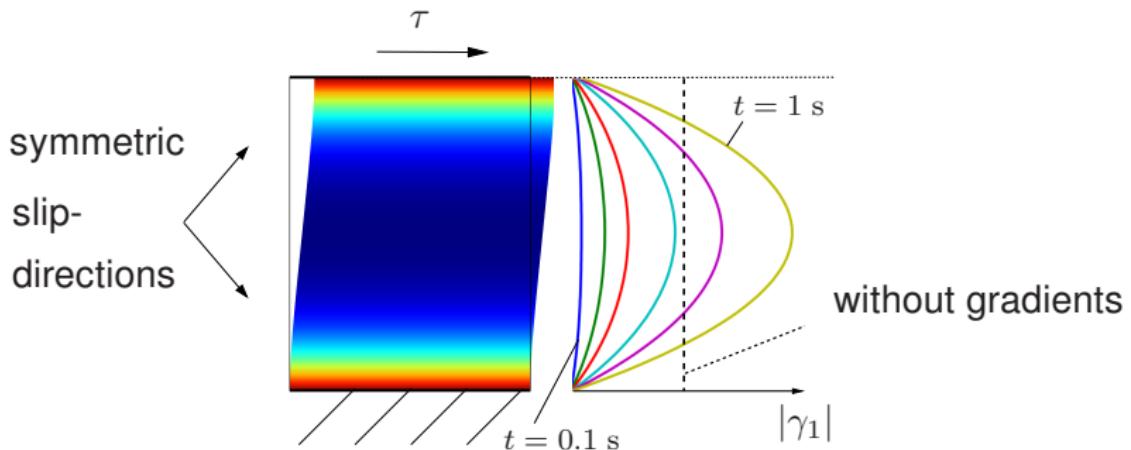


Quadratic defect energy

- leads to smooth strain curves

# Numerical Results

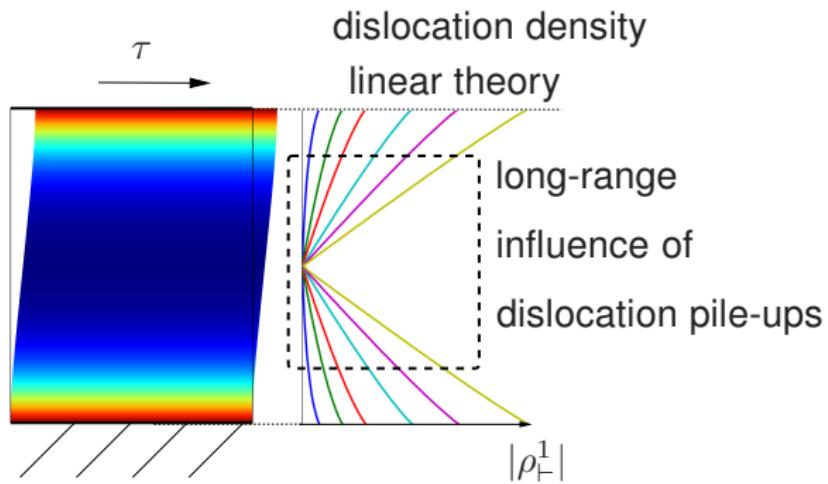
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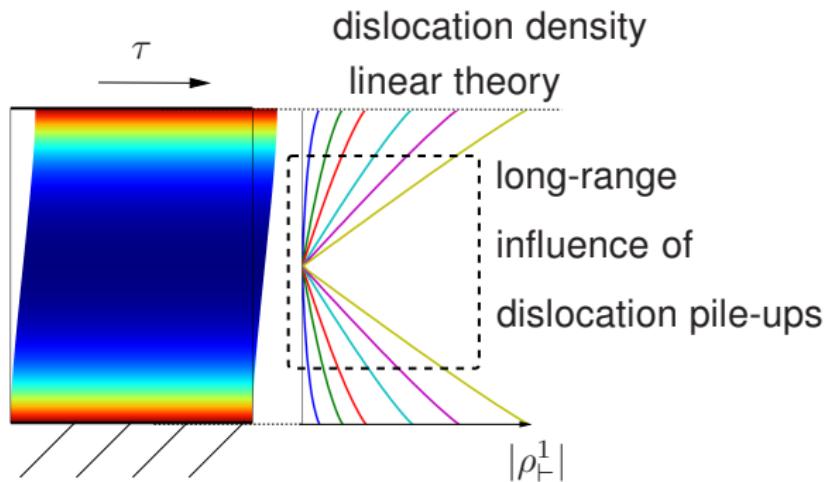
Quadratic defect energy

- leads to smooth strain curves
- long-range dislocation interactions

# Numerical Results



# Numerical Results



Quadratic energy

- leads to a long-range influence of the dislocation pile-ups at the boundaries on the bulk-dislocations

# End of Part I

The support of the German Research Foundation (DFG) in the project "Dislocation based Gradient Plasticity Theory" of the DFG Research Group 1650 "Dislocation based Plasticity" under Grant BO 1466/5-1 is gratefully acknowledged.