

INTRODUCTION

The effective properties of physically non-linear composite materials are investigated using the *nonuniform transformation field analysis (NTFA)* introduced by **Michel and Suquet (2003, 2004)**. An example for the class of examined materials are metal ceramic composites, where ceramic particles are used as a reinforcement embedded into an elasto-plastic matrix material. The influence of *anisotropic particle morphologies* on the physically non-linear effective material response is investigated.

NONUNIFORM TRANSFORMATION FIELD ANALYSIS (Michel and Suquet (2003,2004); Fritzen and Böhlke (2010a))

Finite-dimensional approximation of the *inelastic strain*

$$\boldsymbol{\varepsilon}^p(t, \mathbf{x}) \approx \sum_{\alpha=1}^N \xi_{\alpha}(t) \boldsymbol{\mu}^{(\alpha)}(\mathbf{x})$$

$\boldsymbol{\mu}^{(\alpha)}$: inelastic modes, $\hat{\xi} \in \mathbb{R}^N$: mode activity coefficients

Induced stress and strain fields for given macro-strain $\langle \boldsymbol{\varepsilon} \rangle = \bar{\boldsymbol{\varepsilon}}$

$$\boldsymbol{\varepsilon}(t, \mathbf{x}) = \mathbb{A}(\mathbf{x})[\bar{\boldsymbol{\varepsilon}}(t)] + \sum_{\alpha=1}^N \xi_{\alpha}(t) \boldsymbol{\varepsilon}_*^{(\alpha)}(\mathbf{x})$$

$$\boldsymbol{\sigma}(t, \mathbf{x}) = \mathbb{C}(\mathbf{x})\mathbb{A}(\mathbf{x})[\bar{\boldsymbol{\varepsilon}}(t)] + \sum_{\alpha=1}^N \xi_{\alpha}(t) \boldsymbol{\sigma}_*^{(\alpha)}(\mathbf{x})$$

Thermodynamic driving forces ($\mathbf{B}^{(\alpha)}$: orthonormal basis of Sym)

$$(\hat{A})_{\alpha\beta} = \langle \mathbb{A}^T \boldsymbol{\sigma}_*^{(\alpha)} \cdot \mathbf{B}^{(\beta)} \rangle, \quad (\hat{D})_{\alpha\beta} = \langle \boldsymbol{\mu}^{(\alpha)} \cdot \boldsymbol{\sigma}_*^{(\beta)} \rangle, \quad \hat{\tau} = \hat{A} \hat{\boldsymbol{\varepsilon}} + \hat{D} \hat{\boldsymbol{\xi}}$$

Rate of internal variables

$$\dot{\hat{\boldsymbol{\xi}}} = \dot{\lambda} \frac{\partial \bar{\varphi}}{\partial \hat{\boldsymbol{\tau}}}, \quad \dot{\lambda} \geq 0, \quad \bar{\varphi} \leq 0, \quad \bar{\varphi}(\hat{\boldsymbol{\tau}}) = \|\hat{\boldsymbol{\tau}}\|_2 - k \sigma_F(\bar{q}) \leq 0$$

Effective material response for $\bar{\boldsymbol{\varepsilon}} \equiv \bar{\boldsymbol{\varepsilon}}(t)$, $\hat{\boldsymbol{\xi}} \equiv \hat{\boldsymbol{\xi}}(t)$

$$\bar{\mathbb{C}} = \langle \mathbb{C} \mathbb{A} \rangle, \quad \bar{\boldsymbol{\sigma}}_*^{(\alpha)} = \langle \boldsymbol{\sigma}_*^{(\alpha)} \rangle, \quad \bar{\boldsymbol{\sigma}}(\bar{\boldsymbol{\varepsilon}}, \hat{\boldsymbol{\xi}}) = \bar{\mathbb{C}}[\bar{\boldsymbol{\varepsilon}}] + \sum_{\alpha=1}^N \xi_{\alpha} \bar{\boldsymbol{\sigma}}_*^{(\alpha)}$$

Algorithmic tangent operator

$$\bar{\mathbb{C}}_a = \bar{\mathbb{C}} + \sum_{\alpha=1}^N \langle \boldsymbol{\sigma}_*^{(\alpha)} \rangle \otimes \frac{\partial \xi_{\alpha}}{\partial \boldsymbol{\varepsilon}}$$

ANISOTROPIC MICROSTRUCTURES

(cf. **Fritzen and Böhlke (2011a,b)**)

AIM: Isolate the influence of the particle morphology on the **effective non-linear response**

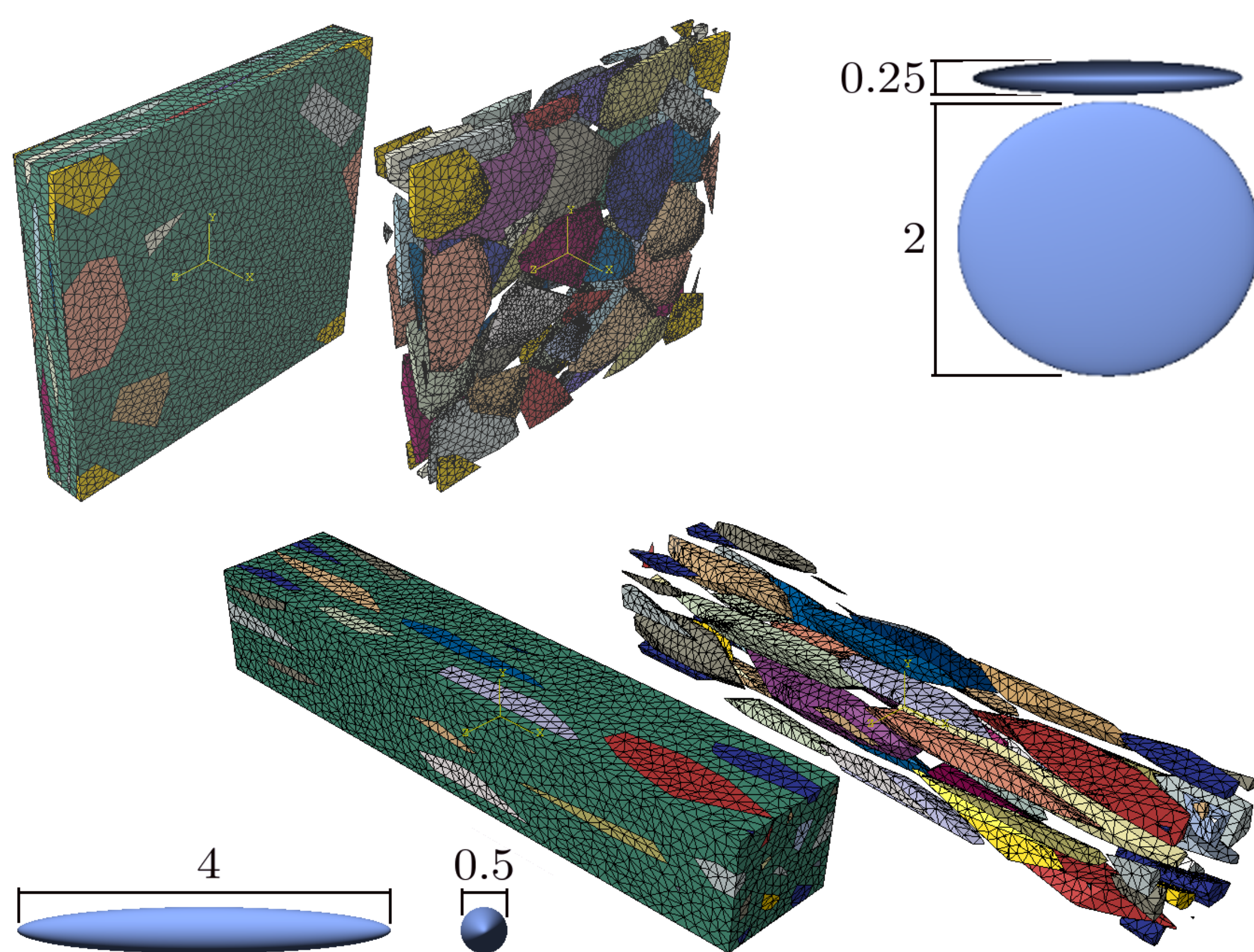


Fig. 1: Anisotropic particle reinforced materials

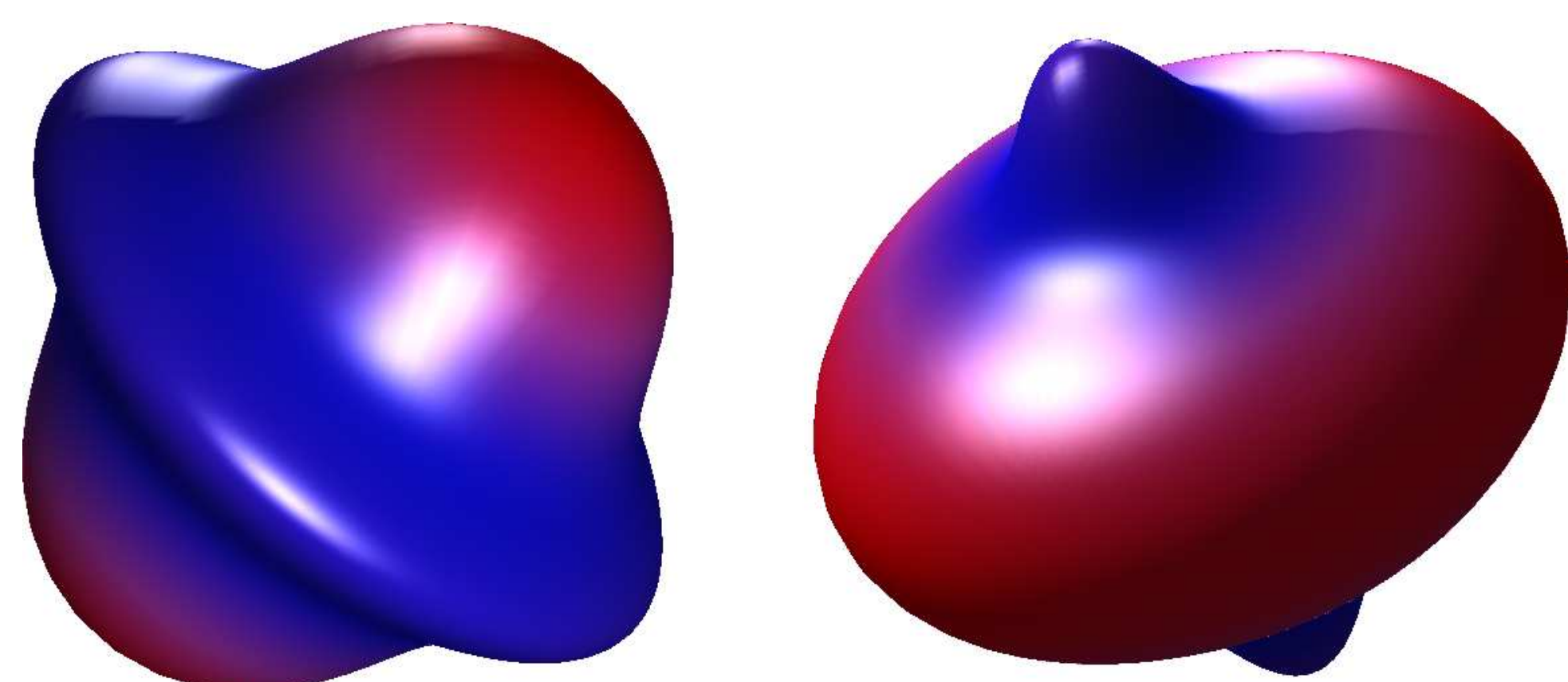


Fig. 2: Effective yield surfaces of the NTFA model

RESULTS (cf. Fritzen and Böhlke (2011b))

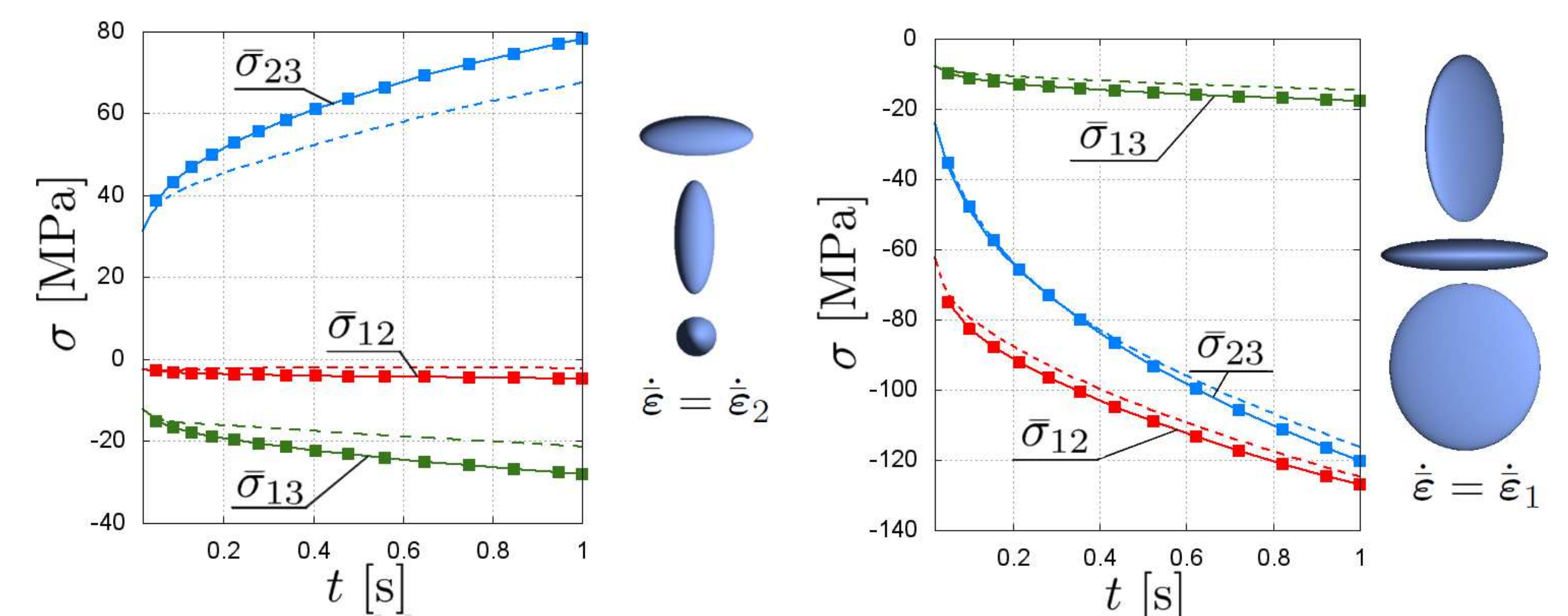


Fig. 3: Comparison of $\bar{\boldsymbol{\sigma}}$: NTFA vs. full-field simulations

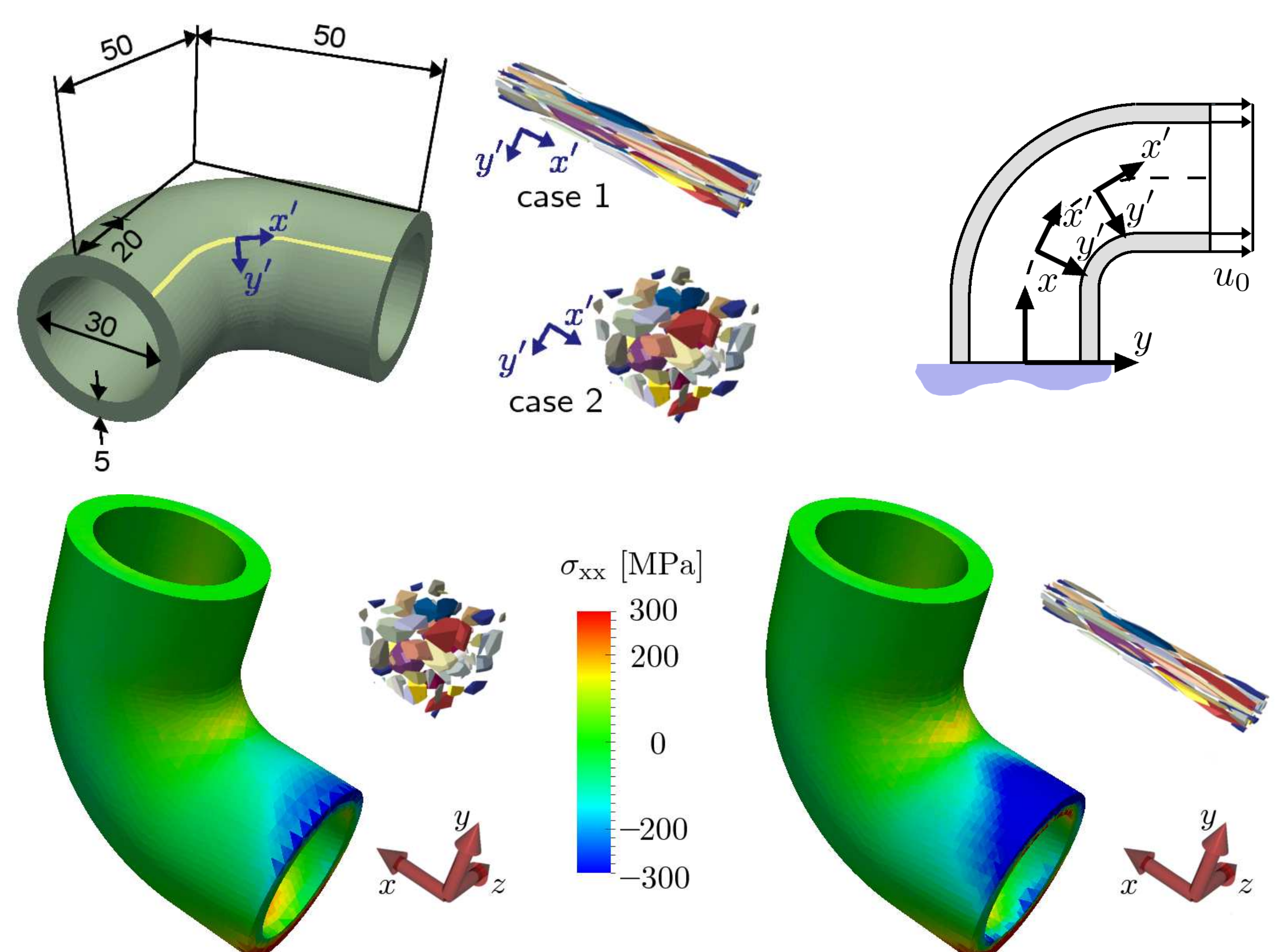


Fig. 4: Influence of the morphology on the macroscopic stress (*left: isotropic particles, right: elongated particles*)

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